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GEOMETRICAL INTERPRETATION OF THE ANTISYMMETRIC COVARIANT DIFFERENTIATION

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Abstract. In this paper it is given a geometrical interpretation of the differentiation (1), where A is a tensor field of type (r,s). Indeed it is proved the equality (5) where $^D_{\phi}$ denotes the parallel displacement along the curve $\tau_4^{-1}\tau_3^{-1}\tau_2^{-1}\tau_1^{-1}$, and $\xi^r = \delta^r_k$, $\eta^S = \delta^S_p$ (fig. 1).

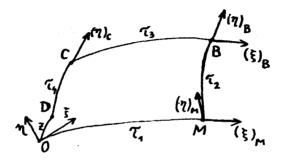


Fig. 1

Let M be an n-dimensional manifold which is endowed with a linear connection. We choose an arbitrary point OEM and a local coordinate system $\mathbf{x}^1,\ldots,\mathbf{x}^n$ in a neighbourhood of O. We shall denote by \mathbf{C}^2 the set of all functions which are twice differentiable and the second derivative is a continuous function in that neighbourhood with respect to the coordinate system (\mathbf{x}^i) .

Let $A_{j_1...j_s}^{i_1...i_r} \in \mathbb{C}^2$ be the tensor field of type (r,s), then

$$(DA)_{j_{1}...j_{s}kp}^{i_{1}...i_{r}} = A_{j_{1}...j_{s};k;p}^{i_{1}...i_{r}} - A_{j_{1}...j_{s};p;k}^{i_{1}...i_{r}}$$
(1)

defines a differentiation which is antisymmetric with respect to the indices k and p. In this paper we will give a geometrical interpretation for this differentiation. Using the Ricci identity, (1) can be written in the following form

$$(DA)_{j_{1}...j_{s}kp}^{i_{1}...i_{r}} = -R^{i_{1}}_{qkp}^{qi_{2}...i_{r}}_{j_{1}...j_{s}} - ... - R^{i_{r}}_{qkp}^{Aj_{1}...j_{s}} + + R^{q}_{j_{1}kp}^{i_{1}...i_{r}}_{qj_{2}...j_{s}} + ... + R^{q}_{j_{s}kp}^{Aj_{1}...q} + T^{q}_{pk}^{Aj_{1}...j_{s};q}.$$

$$(2)$$

Let $\xi, \eta \in T_O(M)$ (fig. 1), and let $\tau_1(t)$ ($0 \le t \le \Delta u$), $\tau_2(t)$ ($0 \le t \le \Delta v$), $\tau_3(t)$ ($0 \le t \le \Delta u$) and $\tau_4(t)$ ($0 \le t \le \Delta v$) be geodesic lines which connect O and M, M and B, B and C, C and D respectively, such that

$$(d\tau_{1}(t)/dt)_{t=0} = \xi,$$
 $(d\tau_{2}(t)/dt)_{t=0} = (\eta)_{M},$ $(d\tau_{3}(t)/dt)_{t=0} = -(\xi)_{B},$ $(d\tau_{4}(t)/dt)_{t=0} = -(\eta)_{C}.$

In the above formulas $(\eta)_M$ is the vector at the point M which is obtained by parallel displacement of η along the curve τ_1 , $(\xi)_B$ is the vector at the point B which is obtained by parallel displacement of ξ along the curve $\tau_1 \cdot \tau_2$, and $(\eta)_C$ is the vector at the point C which is obtained by parallel displacement of η along the curve $\tau_1 \cdot \tau_2 \cdot \tau_3$.

Now we are going to find the differences of the coordinates between the points D and O to the second order of approximation with respect to Δu and $\Delta v.$ Since $\tau_{_{1}}$ is a geodesic line, it holds

$$(x^{k})_{M} - (x^{k})_{O} = (d\tau_{1}^{k}(t)/dt)_{t=o} \Delta u + \frac{1}{2} (d^{2}\tau_{1}^{k}(t)/dt^{2})_{t=o} (\Delta u)^{2} + o(\Delta u^{2}) = \xi^{k} \Delta u - \frac{1}{2} \Gamma_{rs}^{k} \xi^{r} \xi^{s} (\Delta u)^{2} + o(\Delta u^{2}).$$

Similarly we can calculate the differences $(x^k)_B - (x^k)_M$, $(x^k)_C - (x^k)_B$ and $(x^k)_D - (x^k)_C$, and one can obtain that

$$(x^{k})_{D} - (x^{k})_{O} = T_{rs}^{k} \xi^{r} \eta^{s} \Delta u \Delta v + o(\Delta u^{2}) + o(\Delta v^{2}).$$
 (3)

Specially, if $\xi^r = \delta_k^r$ and $\eta^s = \delta_p^s$, (3) implies that

$$(x^{r})_{D} - (x^{r})_{O} = T_{kp}^{r} \Delta u \Delta v + o(\Delta u^{2}) + o(\Delta v^{2}).$$
 (4)

 $\underline{\text{Theorem}}$. Let AEC^2 be a tensor field in a neighbourhood of the point O. Then the following formula holds

$$(DA)_{j_{1}...j_{S}}^{i_{1}...i_{r}}pk = \lim_{\substack{\Delta u \to 0 \\ \Delta v \to 0}} \frac{1}{\Delta u \Delta v} \left[\phi_{O}^{D} (A_{j_{1}...j_{S}}^{i_{1}...i_{r}})_{D} - (A_{j_{1}...j_{S}}^{i_{1}...i_{r}})_{O} \right]$$

$$(5)$$

where $_{\varphi}^{D}$ denotes the parallel displacement along the curve $_{\tau_{4}^{-1}\tau_{3}^{-1}\tau_{2}^{-1}\tau_{1}^{-1}}^{-1}$, and $_{\xi}^{r}=_{\delta}_{k}^{r},~_{\eta}^{s}=_{\delta}_{p}^{s}$ (fig. 1).

<u>Proof.</u> It is well known that if a tensor A at any point of the manifold is transported parallelly round an infinitesimal circuit in the form of an infinitesimal parallelogram with adjacent edges dx^i and δx^i , then the principal part of the infinitesimal increment in $A_{j_1\cdots j_s}^{i_1\cdots i_r}$ is

$$\Delta A_{j_{1} \dots j_{s}}^{i_{1} \dots i_{r}} = (-A_{j_{1} \dots j_{s}}^{q_{1} \dots i_{r}} R_{qkp}^{i_{1}} - \dots - A_{j_{1} \dots j_{s}}^{i_{1} \dots q_{s}} R_{qkp}^{i_{r}} + A_{q_{1} \dots j_{s}}^{i_{1} \dots i_{r}} R_{j_{1} \dots q_{s}}^{q_{1} \dots i_{r}} R_{j_{1} \dots q_{s}}^{q_{1} \dots i_{r}} R_{j_{s} kp}^{q_{s}}) dx^{k} \delta x^{p}.$$
(6)

Let us consider now the closed curve ODCBMO (fig. 1), where the points O and D are connected with a smooth curve z(t). It can be considered as a small distortion of an infinitesimal parallelogram with adjacent edges $\eta^{i} = \delta^{i}_{p}$ and $\xi^{i} = \delta^{i}_{k}$. Let ϕ^{O}_{D} denotes the parallel displacement from O to D along the curve z(t). The formula (6) can be applied and we obtain

$$\phi_{D}^{O}(\phi_{O}^{D}(A_{j_{1}...j_{S}}^{i_{1}...i_{r}})_{D}) - (A_{j_{1}...j_{S}}^{i_{1}...i_{r}})_{D} =$$

$$= -\left[(DA)_{j_{1}...j_{S}}^{i_{1}...i_{r}} + T_{kp}^{q} A_{j_{1}...j_{S}}^{i_{1}...i_{r}} \right] \Delta u \Delta v + o(\Delta u^{2}) + o(\Delta v^{2}).$$

Applying the formula for the parallel displacement of the tensor ${}^D_0(A_{j_1...j_S}^{i_1...i_r})$ from 0 to D, and using that

$$(A_{j_1,\ldots,j_s}^{i_1,\ldots,i_r})_D = (A_{j_1,\ldots,j_s}^{i_1,\ldots,i_r})_O + \partial (A_{j_1,\ldots,j_s}^{i_1,\ldots,i_r})/\partial x^k ((x^k)_D - (x^k)_O) + O(\Delta u^2) + O(\Delta v^2)$$

we obtain

$$\phi_{0}^{D}(A_{j_{1}...j_{s}}^{i_{1}...i_{r}})_{D} - (A_{j_{1}...j_{s}}^{i_{1}...i_{r}})_{O} = (DA)_{j_{1}...j_{s}}^{i_{1}...i_{r}} \psi_{0}^{\Delta u \Delta v} + o(\Delta u^{2}) + o(\Delta v^{2}).$$

If we devide this equality by $\Delta u \Delta v$ and put $\Delta u \rightarrow 0$, $\Delta v \rightarrow 0$ we will obtain the formula (5).

The above theorem gives the required geometrical interpretation of the antisymmetric differentiation D, using the parallel displacement of tensors. It is well known that the covariant differentiation can be written in the following form

$$A_{j_{1}...j_{s};k}^{i_{1}...i_{r}} = \lim_{\Delta u \to o} \frac{1}{\Delta u} \left[\phi_{M}^{O}(A_{j_{1}...j_{s}}^{i_{1}...i_{r}}) - (A_{j_{1}...j_{s}}^{i_{1}...i_{r}})_{M} \right],$$

where ϕ_{M}^{O} denotes the parallel displacement from O to M along the geodesic line $\tau_{1}(t)$ (O $\leq t \leq \Delta u$) (fig. 1). We notice that both of these differentiations have similar geometrical interpretations.

Similar result can be obtained if one does not use indices. Indeed, if AEC^2 is a tensor field in a neighbourhood of the point O and $X,YET_O(M)$, then the following formula

$$D_{Y,X}(A) = \lim_{\substack{\Delta u \to o \\ \Delta v \to o}} \frac{1}{\Delta u \Delta v} \left[\phi_O^D(A_D) - A_O \right]$$
 (7)

holds, where

$$D_{\mathbf{Y},\mathbf{X}} = \nabla_{\mathbf{X}} \cdot \nabla_{\mathbf{Y}} - \nabla_{\mathbf{Y}} \cdot \nabla_{\mathbf{X}} + \nabla_{(\nabla_{\mathbf{Y}}\mathbf{X})} - \nabla_{(\nabla_{\mathbf{X}}\mathbf{Y})}.$$

We note that $\Delta u/\Delta v$ and $\Delta v/\Delta u$ in (5) and (7) are supposed to be bounded numbers.

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ГЕОМЕТРИСКА ИНТЕРПРЕТАЦИЈА НА АНТИСИМЕТРИЧНОТО КОВАРИЈАНТНО ДИФЕРЕНЦИРАЊЕ

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Резиме

Во овој труд е дадена геометриска интерпретација на диференцирањето дефинирано со (1) односно (2), каде што A е тензорно поле од тип (r,s). Всушност, се докажува равенството (5) каде што ф означува паралелно пренесување по должината на кривата τ_4^{-1} τ_3^{-1} τ_4^{-1} , и $\xi^r = \delta_k^r$, $\eta^s = \delta_p^s$ (сл. 1).